# Banded symmetric Toeplitz matrices: where linear algebra borrows from difference equations 

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A Toeplitz matrix, named after the German mathematician Otto Toeplitz (1881-1940), is of the form $T=\left[t_{r-s}\right]_{r, s=0}^{n-1}$. (It's ok, and convenient for Toeplitz matrices, to number rows and columns from 0 to $n-1$.) A symmetric Toeplitz matrix is of the form $T_{n}=\left[t_{|r-s|}\right]_{r, s=0}^{n-1}$. For example,

$$
T_{5}=\left[\begin{array}{ccccc}
t_{0} & t_{1} & t_{2} & t_{3} & t_{4} \\
t_{1} & t_{0} & t_{1} & t_{2} & t_{3} \\
t_{2} & t_{1} & t_{0} & t_{1} & t_{2} \\
t_{3} & t_{2} & t_{1} & t_{0} & t_{1} \\
t_{4} & t_{3} & t_{2} & t_{1} & t_{0}
\end{array}\right]
$$

is a $5 \times 5$ symmetric Toeplitz matrix. We will assume that $t_{0}, t_{1}, \ldots, t_{k-1}$ are all real numbers. From your linear algebra course you know that a symmetric matrix with real entries has real eigenvalues and is always diagonalizable; that is, $T_{n}$ has real eigenvalues and $n$ linearly independent eigenvectors.

A Toeplitz matrix is said to be banded if there is an integer $d<n-1$ such that $t_{\ell}=0$ if $\ell>d$. In this case, we say that $T$ has bandwidth $d$. For example,

$$
T_{5}=\left[\begin{array}{ccccc}
t_{0} & t_{1} & t_{2} & 0 & 0 \\
t_{1} & t_{0} & t_{1} & t_{2} & 0 \\
t_{2} & t_{1} & t_{0} & t_{1} & t_{2} \\
0 & t_{2} & t_{1} & t_{0} & t_{1} \\
0 & 0 & t_{2} & t_{1} & t_{0}
\end{array}\right]
$$

is a $5 \times 5$ banded symmetric Toeplitz matrix with bandwidth 2 .

The eigenvalue problem for very large ( $n$ can be in the thousands!) symmetric banded Toeplitz matrices pops up in many statistical problems. In your linear algebra course you learned to solve the eigenvalue problem for a matrix $A$ by factoring its characteristic polynomial

$$
p(\lambda)=\operatorname{det}(A-\lambda I) .
$$

Sorry, that's impossible for big matrices. In general there is no computationally useful way to obtain the characteristic polynomial of a large symmetric matrix (or any other large matrix). All methods for finding a single eigenvalue of an arbitrary $n \times n$ symmetric matrix carry a computational cost (it's called complexity) proportional to $n^{3}$. So, if you double the size of the matrix you make the problem of obtaining a single eigenvalue eight times more difficult. However, the situation is different for banded symmetric Toeplitz matrices.

Let's start with the simplest case: $d=1$.

$$
T_{n}=\left[\begin{array}{rrrrrrrr}
t_{0} & t_{1} & 0 & 0 & \cdots & 0 & 0 & 0 \\
t_{1} & t_{0} & t_{1} & 0 & \cdots & 0 & 0 & 0 \\
0 & t_{1} & t_{0} & t_{1} & \cdots & 0 & 0 & 0 \\
\vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\
0 & 0 & 0 & 0 & \cdots & t_{1} & t_{0} & t_{1} \\
0 & 0 & 0 & 0 & \cdots & 0 & t_{1} & t_{0}
\end{array}\right]_{n \times n}
$$

with $t_{1} \neq 0$. This is a symmetric tridiagonal Toeplitz matrix. A vector

$$
x=\left[\begin{array}{c}
x_{0} \\
x_{1} \\
\vdots \\
x_{n-1}
\end{array}\right]
$$

is a $\lambda$-eigenvector of $T_{n}$ if and only if

$$
\begin{aligned}
t_{0} x_{0}+t_{1} x_{1} & =\lambda x_{0} \\
t_{1} x_{j-1}+t_{1} x_{0}+t_{1} x_{j+1} & =\lambda x_{j}, \quad 1 \leq j \leq n-1, \\
t_{1} x_{n-1}+t_{0} x_{n} & =\lambda x_{n}
\end{aligned}
$$

which we can rewrite as

$$
t_{1} x_{j-1}+\left(t_{0}-\lambda\right) x_{j}+t_{1} x_{j+1}=0, \quad 0 \leq j \leq n,
$$

(a homogeneous difference equation) if we define $x_{0}=$ $x_{n+1}=0$ (boundary conditions).

The characteristic polynomial of the difference equation

$$
\begin{equation*}
t_{1} x_{j-1}+\left(t_{0}-\lambda\right) x_{j}+t_{1} x_{j+1}=0 \tag{DE}
\end{equation*}
$$

is
$p(z ; \lambda)=t_{1}+\left(t_{0}-\lambda\right) z+t_{1} z^{2}=t_{1}\left(z-z_{1}(\lambda)\right)\left(z-z_{2}(\lambda)\right) ;$ thus, $p\left(z_{1}(\lambda)\right)=p\left(z_{2}(\lambda)\right)=0$. (We don't know $z_{1}(\lambda)$ and $z_{2}(\lambda)$ yet; be patient.) If we let

$$
x_{j}=c_{1} z_{1}^{j}(\lambda)+c_{2} z_{2}^{j}(\lambda)
$$

where $c_{1}$ and $c_{2}$ are arbitrary constants, then the left side of (DE) equals

$$
c_{1} z_{1}^{j-1} p\left(z_{1}(\lambda)\right)+c_{2} z_{2}^{j-1} p\left(z_{2}(\lambda)\right)=0
$$

for any choice of $c_{1}$ and $c_{2}$. Now let's work on the boundary conditions. Since $x_{0}=0$ if and only if $c_{2}=-c_{1}$,

$$
x_{j}=c\left(z_{1}^{j}(\lambda)-z_{2}^{j}(\lambda)\right)
$$

Now $x_{n+1}=0$ if and only if $\left(z_{1}(\lambda) / z_{2}(\lambda)\right)^{n+1}=1$, which is true if and only if
$z_{1}(\lambda)=\gamma_{q} \exp \left(\frac{q \pi i}{n+1}\right) \quad$ and $\quad z_{2}(\lambda)=\gamma_{q} \exp \left(\frac{-q \pi i}{n+1}\right)$,
where $\exp (i \theta)=e^{i \theta}=\cos \theta+i \sin \theta, q=1, \ldots, n$ and $\gamma_{q}$ is to be determined. (Letting $q=0$ does not produce an eigenvector because if $\left.z_{1}(\lambda)=z_{2}(\lambda)\right)$ then $x_{j}=0$ for all $j)$.

Taking note that are $q$ possibilities, the eigenvectors have the form

$$
x_{q}=\left[\begin{array}{c}
x_{0 q} \\
x_{1 q} \\
\vdots \\
x_{n-1, q}
\end{array}\right]
$$

where

$$
x_{j q}=c\left(z_{1}^{j}\left(\lambda_{q}\right)-z_{2}^{j}\left(\lambda_{q}\right)\right)
$$

$z_{1 q}(\lambda)=\gamma_{q} \exp \left(\frac{q \pi i}{n+1}\right), \quad$ and $\quad z_{2 q}(\lambda)=\gamma_{q} \exp \left(\frac{-q \pi i}{n+1}\right)$,
so

$$
x_{j q}=c\left(\exp \frac{j q \pi i}{n+1}-\exp \frac{j q \pi i}{n+1}\right)=2 c i \sin \frac{j q \pi}{n+1} .
$$

Since $c$ is arbitrary, it makes sense to let $c=1 / 2 \gamma_{q} i$. (Don't worry that maybe $\gamma_{q}=0$; we'll see that it isn't.) Then

$$
x_{j q}=\sin \frac{j q \pi}{n+1}, \quad 0 \leq j \leq n-1
$$

ALL SYMMETRIC TRIDIAGONAL TOEPLITZ MATRICES HAVE THE SAME EIGENVECTORS!

Now let's find $\lambda_{q}$, the eigenvalue associated with $q$.

$$
t_{1}+\left(t_{0}-\lambda_{q}\right)+t_{1} z^{2}=t_{1}\left(z-z_{1}(\lambda)\right)\left(z_{2}-z_{2}(\lambda)\right)
$$

which equals

$$
t_{1}\left(z^{2}-\left(z_{1}(\lambda)+z_{2}(\lambda)\right) z+z_{1}(\lambda) z_{2}(\lambda)\right)
$$

Since

$$
\begin{aligned}
& z_{1}(\lambda)=\gamma_{q} \exp \left(\frac{q \pi i}{n+1}\right) \quad \text { and } \quad z_{2}(\lambda)=\gamma_{q} \exp \left(\frac{-q \pi i}{n+1}\right) \\
& t_{1}+\left(t_{0}-\lambda_{q}\right) z+t_{1} z^{2}=t_{1}\left(z^{2}-2 \gamma_{q} z \cos \left(\frac{q \pi}{n+1}\right)+\gamma_{q}^{2}\right)
\end{aligned}
$$

Equating coefficients on the two sides yields $\gamma_{q}=1$ and

$$
\lambda_{q}=t_{0}+2 t_{1} \cos \left(\frac{q \pi}{n+1}\right), \quad 1 \leq q \leq n
$$

Now suppose $d>1$. To see where we're going, a nonzero vector

$$
x=\left[\begin{array}{l}
x_{0} \\
x_{1} \\
x_{2} \\
x_{3} \\
x_{4}
\end{array}\right]
$$

is a $\lambda$-eigenvector of

$$
T_{5}=\left[\begin{array}{ccccc}
t_{0} & t_{1} & t_{2} & 0 & 0 \\
t_{1} & t_{0} & t_{1} & t_{2} & 0 \\
t_{2} & t_{1} & t_{0} & t_{1} & t_{2} \\
0 & t_{2} & t_{1} & t_{0} & t_{1} \\
0 & 0 & t_{2} & t_{1} & t_{0}
\end{array}\right]
$$

if and only if

$$
\left[\begin{array}{ccccc}
t_{0} & t_{1} & t_{2} & 0 & 0 \\
t_{1} & t_{0} & t_{1} & t_{2} & 0 \\
t_{2} & t_{1} & t_{0} & t_{1} & t_{2} \\
0 & t_{2} & t_{1} & t_{0} & t_{1} \\
0 & 0 & t_{2} & t_{1} & t_{0}
\end{array}\right]\left[\begin{array}{l}
x_{0} \\
x_{1} \\
x_{2} \\
x_{3} \\
x_{4}
\end{array}\right]=\lambda\left[\begin{array}{l}
x_{0} \\
x_{1} \\
x_{2} \\
x_{3} \\
x_{4}
\end{array}\right]
$$ or, equivalently,

$$
\begin{aligned}
t_{0} x_{0}+t_{1} x_{1}+t_{2} x_{2} & =\lambda x_{0} \\
t_{1} x_{0}+t_{0} x_{1}+t_{1} x_{2}+t_{2} x_{3} & =\lambda x_{1} \\
t_{2} x_{0}+t_{1} x_{1}+t_{2} x_{2}+t_{3} x_{3}+t_{4} x_{4} & =\lambda x_{2} \\
t_{2} x_{1}+t_{1} x_{2}+t_{0} x_{3}+t_{1} x_{4} & =\lambda x_{3} \\
t_{2} x_{2}+t_{1} x_{3}+t_{0} x_{4} & =\lambda x_{4}
\end{aligned}
$$

(repeated for clarity)

$$
\begin{aligned}
t_{0} x_{0}+t_{1} x_{1}+t_{2} x_{2} & =\lambda x_{0} \\
t_{1} x_{0}+t_{0} x_{1}+t_{1} x_{2}+t_{2} x_{3} & =\lambda x_{1} \\
t_{2} x_{0}+t_{1} x_{1}+t_{2} x_{2}+t_{3} x_{3}+t_{4} x_{4} & =\lambda x_{2} \\
t_{2} x_{1}+t_{1} x_{2}+t_{0} x_{3}+t_{1} x_{4} & =\lambda x_{3} \\
t_{2} x_{2}+t_{1} x_{3}+t_{0} x_{4} & =\lambda x_{4}
\end{aligned}
$$

or, equivalently,

$$
\begin{aligned}
t_{2} x_{-2}+t_{1} x_{-1}+t_{0} x_{0}+t_{1} x_{1}+t_{2} x_{2} & =\lambda x_{0} \\
t_{2} x_{-1}+t_{1} x_{0}+t_{0} x_{1}+t_{1} x_{2}+t_{2} x_{3} & =\lambda x_{1} \\
t_{2} x_{0}+t_{1} x_{1}+t_{0} x_{2}+t_{1} x_{3}+t_{2} x_{4} & =\lambda x_{2} \\
t_{2} x_{1}+t_{1} x_{2}+t_{0} x_{3}+t_{1} x_{4}+t_{2} x_{5} & =\lambda x_{3} \\
t_{2} x_{2}+t_{1} x_{3}+t_{0} x_{4}+t_{1} x_{5}+t_{2} x_{6} & =\lambda x_{4}
\end{aligned}
$$

if we impose the boundary conditions

$$
x_{-2}=x_{-1}=x_{5}=x_{6}=0
$$

Better yet,

$$
\sum_{\ell=-2}^{2} t_{|\ell|} x_{\ell+r}=\lambda x_{r}, \quad 0 \leq r \leq 4
$$

(repeated for clarity)

$$
\sum_{\ell=-2}^{2} t_{|\ell|} x_{\ell+r}=\lambda x_{r}, \quad 0 \leq r \leq 4
$$

with boundary conditions

$$
x_{-2}=x_{-1}=x_{5}=x_{6}=0
$$

For the general case where $T_{n}=\left[t_{|r-s|}\right]_{r, s=0}^{n-1}$ with $t_{\ell}=0$ if $\ell>d$, the eigenvalue problem can be written as

$$
\begin{equation*}
\sum_{\ell=-d}^{d} t_{|\ell|} x_{\ell+r}=\lambda x_{r}, \quad 0 \leq r \leq n-1 \tag{DE}
\end{equation*}
$$

subject to

$$
x_{r}=0, \quad-d \leq r \leq-1, \quad n \leq r \leq n+d-1
$$

Eqn. (DE) is a difference equation and the conditions in (BC) are called boundary conditions. Obviously, (DE) and (BC) both hold for any $\lambda$ if $x_{r}=0$ for $-d \leq r \leq n+d-1$. However, that's not interesting, since an eigenvector must be nonzero. Finding the values of $\lambda$ for which (DE) has nonzero solutions that satisfy (BC) is a boundary value problem.

The characteristic polynomial of the difference equation

$$
\sum_{\ell=-d}^{d} t_{|\ell|} x_{\ell+r}=\lambda x_{r}, \quad 0 \leq r \leq n-1
$$

is

$$
P(z, \lambda)=\sum_{\ell=-d}^{d} t_{|\ell|} z^{\ell}-\lambda
$$

The zeros of $P(z, \lambda)$ are continuous functions of $\lambda$ and, since $P(z, \lambda)=P(1 / z, \lambda)$, they occur in reciprocal pairs

$$
\left(z_{1}(\lambda), 1 / z_{1}(\lambda)\right), \ldots,\left(z_{d}(\lambda), 1 / z_{d}(\lambda)\right)
$$

It can be shown (don't you hate that?) that these zeros are distinct except for at most finitely many "bad values" of $\lambda$. We'll assume that none of these bad values are actually eigenvalues of $T_{n}$. (This is a pretty safe bet.) Then (DE) holds if
$x_{r}=\sum_{s=1}^{d}\left(a_{s} z_{s}^{r}(\lambda)+b_{s} z_{s}^{-r}(\lambda)\right), \quad-d \leq r \leq n+d-1$,
where $a_{1}, \ldots, a_{d}$ and $b_{1}, \ldots, b_{d}$ are arbitrary constants.

Proof. Recall that $P(z, \lambda)=\sum_{\ell=-d}^{d} t_{|\ell|} z^{\ell}-\lambda$. If

$$
x_{r}=\sum_{s=1}^{d}\left(a_{s} z_{s}^{r}(\lambda)+b_{s} z_{s}^{-r}(\lambda)\right),-d \leq r \leq n+d-1, \text { then }
$$

$$
\begin{aligned}
\sum_{\ell=-d}^{d} t_{|\ell|} x_{\ell+r}-\lambda x_{r}= & \sum_{\ell=-d}^{d} t_{|\ell|} \sum_{s=1}^{d}\left(a_{s} z_{s}^{\ell+r}+b_{s} z_{s}^{-\ell-r}\right) \\
& -\lambda \sum_{s=1}^{d}\left(a_{s} z_{s}^{r}(\lambda)+b_{s} z_{s}^{-r}(\lambda)\right) \\
= & \sum_{s=1}^{d} a_{s} z_{s}^{r}(\lambda)\left(\sum_{\ell=-d}^{d} t_{|\ell|} z_{s}^{\ell}(\lambda)-\lambda\right)
\end{aligned}
$$

$$
+\sum_{s=1}^{d} b_{s} z_{s}^{-r}(\lambda)\left(\sum_{\ell=-d}^{d} t_{|\ell|} z_{s}^{-\ell}(\lambda)-\lambda\right)
$$

$$
=\sum_{s=1}^{d}\left(a_{s} z_{s}^{r}(\lambda) P\left(z_{s}(\lambda), \lambda\right)+b_{s} z_{s}^{-r}(\lambda) P\left(1 / z_{s}(\lambda), \lambda\right)\right)=0
$$

(look at the top of this page) for all $r$. Note that $a_{1}, \ldots, a_{d}$ and $b_{1}, \ldots, b_{d}$ are completely arbitrary up to this point.

Now we must choose them to satisfy the boundary conditions

$$
x_{r}=0, \quad-d \leq r \leq-1, \quad n \leq r \leq n+d-1 ; \quad \text { (BC) }
$$

that is,

$$
\sum_{s=1}^{d}\left(a_{s} z_{s}^{r}(\lambda)+b_{s} z_{s}^{-r}(\lambda)\right)=0
$$

for $-d \leq r \leq-1$ and $n \leq r \leq n+d-1$. For example, if $d=2$, we must have

$$
\begin{aligned}
& {\left[\begin{array}{cccc}
z_{1}^{-1}(\lambda) & z_{2}^{-1}(\lambda) & z_{1}(\lambda) & z_{2}(\lambda) \\
z_{1}^{2}(\lambda) & z_{2}^{-2}(\lambda) & z_{1}^{2}(\lambda) & z_{2}^{2}(\lambda) \\
z_{1}^{n}(\lambda) & z_{2}^{n}(\lambda) & z_{1}^{-n}(\lambda) & z_{2}^{-n}(\lambda) \\
z_{1}^{n+1}(\lambda) & z_{2}^{n+1}(\lambda) & z_{1}^{-n-1}(\lambda) & z_{2}^{-n-1}(\lambda)
\end{array}\right]\left[\begin{array}{l}
a_{1} \\
a_{2} \\
b_{1} \\
b_{2}
\end{array}\right] } \\
= & {\left[\begin{array}{l}
0 \\
0 \\
0 \\
0
\end{array}\right] }
\end{aligned}
$$

For clarity,

$$
\left[\begin{array}{cccc}
z_{1}^{-1}(\lambda) & z_{2}^{-1}(\lambda) & z_{1}(\lambda) & z_{2}(\lambda) \\
z_{1}^{-2}(\lambda) & z_{2}^{-2}(\lambda) & z_{1}^{2}(\lambda) & z_{2}^{2}(\lambda) \\
z_{1}^{n}(\lambda) & z_{2}^{n}(\lambda) & z_{1}^{-n}(\lambda) & z_{2}^{-n}(\lambda) \\
z_{1}^{n+1}(\lambda) & z_{2}^{n+1}(\lambda) & z_{1}^{-n-1}(\lambda) & z_{2}^{-n-1}(\lambda)
\end{array}\right]\left[\begin{array}{l}
a_{1} \\
a_{2} \\
b_{1} \\
b_{2}
\end{array}\right]=0
$$

Let

$$
\begin{gathered}
P(\lambda)=\left[\begin{array}{cc}
z_{1}^{-1} \lambda & z_{2}^{-1}(\lambda) \\
z_{1}^{-2} \lambda & z_{2}^{-2}(\lambda)
\end{array}\right], \quad Q(\lambda)=\left[\begin{array}{cc}
z_{1} \lambda & z_{2}(\lambda) \\
z_{1}^{2} \lambda & z_{2}^{2}(\lambda)
\end{array}\right] \\
R_{n}(\lambda)=\left[\begin{array}{cc}
z_{1}^{n} \lambda & z_{2}^{n}(\lambda) \\
z_{1}^{n+1} \lambda & z_{2}^{n+1}(\lambda)
\end{array}\right] \\
S_{n}(\lambda)=\left[\begin{array}{cc}
z_{1}^{-n} \lambda & z_{2}^{-n}(\lambda) \\
z_{2}^{-n-1} \lambda & z_{2}^{-n-1}(\lambda)
\end{array}\right] \\
a=\left[\begin{array}{c}
a_{1} \\
a_{2}
\end{array}\right], \quad b=\left[\begin{array}{c}
b_{1} \\
b_{2}
\end{array}\right]
\end{gathered}
$$

Then the boundary conditions are satisfied if and only if

$$
\left[\begin{array}{cc}
P(\lambda) & Q(\lambda) \\
R_{n}(\lambda) & S_{n}(\lambda)
\end{array}\right]\left[\begin{array}{l}
a \\
b
\end{array}\right]=0
$$

In general, let

$$
\begin{aligned}
P(\lambda) & =\left[z_{s}^{-r}(\lambda)\right]_{r, s=1}^{d}, \quad Q(\lambda)=\left[z_{s}^{r}(\lambda)\right], \\
R_{n}(\lambda) & =\left[z_{s}^{n+r-1}(\lambda)\right], \quad S_{n}(\lambda)=\left[z_{s}^{-n+r-1}(\lambda)\right] \\
a & =\left[\begin{array}{c}
a_{1} \\
a_{2} \\
\vdots \\
a_{d}
\end{array}\right] \quad \text { and } \quad b=\left[\begin{array}{c}
b_{1} \\
b_{2} \\
\vdots \\
b_{d}
\end{array}\right] .
\end{aligned}
$$

Then the boundary conditions are satisfied if and only if

$$
\left[\begin{array}{cc}
P(\lambda) & Q(\lambda)  \tag{S}\\
R_{n}(\lambda) & S_{n}(\lambda)
\end{array}\right]\left[\begin{array}{l}
a \\
b
\end{array}\right]=0 .
$$

Let

$$
D_{n}(\lambda)=\left|\begin{array}{cc}
P(\lambda) & Q(\lambda) \\
R_{n}(\lambda) & S_{n}(\lambda)
\end{array}\right| \quad \text { (determinant). }
$$

An eigenvector of $T_{n}$ must be a nonzero vector. Since (S) has only the trivial solution $\left[\begin{array}{l}a \\ b\end{array}\right]=\left[\begin{array}{l}0 \\ 0\end{array}\right]$ if $D_{n}(\lambda) \neq$ 0 , it follows that $\lambda$ is an eigenvalue of $T_{n}$ if and only if $D_{n}(\lambda)=0$. For ways to find the zeros of $D_{n}(\lambda)$, see my papers RP-44, 61, 63, and 78. Since $D_{n}(\lambda)$ doesn't become more complicated as $n$ increases, the difficulty of finding individual eigenvalues of $T_{n}$ is independent of $n$.

We can take this a little further. The eigenvectors of a symmetric Toeplitz matrix have a special property that I haven't mentioned. To identify this property, let $J_{n}$ be the "flip matrix," which has 1's on its secondary diagonal and 0's elsewhere. For example,

$$
J_{5}=\left[\begin{array}{lllll}
0 & 0 & 0 & 0 & 1 \\
0 & 0 & 0 & 1 & 0 \\
0 & 0 & 1 & 0 & 0 \\
0 & 1 & 0 & 0 & 0 \\
1 & 0 & 0 & 0 & 0
\end{array}\right]
$$

Note that

$$
\begin{aligned}
J_{5}^{2} & =\left[\begin{array}{lllll}
0 & 0 & 0 & 0 & 1 \\
0 & 0 & 0 & 1 & 0 \\
0 & 0 & 1 & 0 & 0 \\
0 & 1 & 0 & 0 & 0 \\
1 & 0 & 0 & 0 & 0
\end{array}\right]\left[\begin{array}{lllll}
0 & 0 & 0 & 0 & 1 \\
0 & 0 & 0 & 1 & 0 \\
0 & 0 & 1 & 0 & 0 \\
0 & 1 & 0 & 0 & 0 \\
1 & 0 & 0 & 0 & 0
\end{array}\right] \\
& =\left[\begin{array}{lllll}
1 & 0 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 & 0 \\
0 & 0 & 1 & 0 & 0 \\
0 & 0 & 0 & 1 & 0 \\
0 & 0 & 0 & 0 & 1
\end{array}\right]=I_{5}
\end{aligned}
$$

In general, $J_{n}^{2}=I_{n}$; that is, $J_{n}$ is its own inverse.

Multiplying a vector by $J_{n}$ reverses ("flips") the components of the vector. For example, if

$$
x=\left[\begin{array}{l}
x_{0} \\
x_{1} \\
x_{2} \\
x_{3} \\
x_{4}
\end{array}\right]
$$

then

$$
J_{5} x=\left[\begin{array}{ccccc}
0 & 0 & 0 & 0 & 1 \\
0 & 0 & 0 & 1 & 0 \\
0 & 0 & 1 & 0 & 0 \\
0 & 1 & 0 & 0 & 0 \\
1 & 0 & 0 & 0 & 0
\end{array}\right]\left[\begin{array}{l}
x_{0} \\
x_{1} \\
x_{2} \\
x_{3} \\
x_{4}
\end{array}\right]=\left[\begin{array}{l}
x_{4} \\
x_{3} \\
x_{2} \\
x_{1} \\
x_{0}
\end{array}\right]
$$

In general, if

$$
x=\left[\begin{array}{c}
x_{0} \\
x_{1} \\
\vdots \\
x_{n-2} \\
x_{n-1}
\end{array}\right] \text { then } J_{n} x=\left[\begin{array}{c}
x_{n-1} \\
x_{n-2} \\
\vdots \\
x_{1} \\
x_{0}
\end{array}\right]
$$

We say that a vector $x$ is symmetric if $J_{n} x=x$, or skewsymmetric if $J_{n} x=-x$. Adding symmetric vectors produces a symmetric vector, and multiplying a symmetric vector by a real number produces a symmetric vector; hence, the symmetric vectors in $\mathbb{R}^{n}$ form a subspace of $\mathbb{R}^{n}$. Similarly, the skew-symmetric vectors form a subspace of $\mathbb{R}^{n}$. If $n=2 m$ then each of these subspaces has dimension $m$. If $n=2 m+1$ then the subspace of symmetric vectors has dimension $m+1$ and the subspace of skew symmetric vectors has dimension $m$. The zero vector is the only vector that is both symmetric and skew-symmetric.

For example, if $n=4$ then

$$
\left[\begin{array}{c}
1 \\
0 \\
0 \\
1
\end{array}\right] \text { and }\left[\begin{array}{l}
0 \\
1 \\
1 \\
0
\end{array}\right]
$$

form a basis for the subspace of symmetric vectors, while

$$
\left[\begin{array}{r}
1 \\
0 \\
0 \\
-1
\end{array}\right] \text { and }\left[\begin{array}{r}
0 \\
1 \\
-1 \\
0
\end{array}\right]
$$

form a basis for the subspace of skew-symmetric vectors.

If $n=5$ then

$$
\left[\begin{array}{l}
1 \\
0 \\
0 \\
0 \\
1
\end{array}\right], \quad\left[\begin{array}{l}
0 \\
1 \\
0 \\
1 \\
0
\end{array}\right], \quad \text { and } \quad\left[\begin{array}{c}
0 \\
0 \\
1 \\
0 \\
0
\end{array}\right]
$$

form a basis for the subspace of symmetric vectors, while

$$
\left[\begin{array}{r}
1 \\
0 \\
0 \\
0 \\
-1
\end{array}\right] \text { and }\left[\begin{array}{r}
0 \\
1 \\
0 \\
-1 \\
0
\end{array}\right]
$$

form a basis for the subspace of skew-symmetric vectors. In general, if $n=2 m$ then the subspace of symmetric vectors and the subspace of skew-symmetric vectors are both $m$-dimensional. If $n=2 m+1$ then the subspace of symmetric vectors is $(m+1)$-dimensional and the subspace of skew-symmetric vectors is $m$-dimensional.

Multiplying a matrix on the left by $J_{n}$ reverses the rows of the matrix, so

$$
\begin{aligned}
J_{5} T_{5} & =\left[\begin{array}{lllll}
0 & 0 & 0 & 0 & 1 \\
0 & 0 & 0 & 1 & 0 \\
0 & 0 & 1 & 0 & 0 \\
0 & 1 & 0 & 0 & 0 \\
1 & 0 & 0 & 0 & 0
\end{array}\right]\left[\begin{array}{lllll}
t_{0} & t_{1} & t_{2} & t_{3} & t_{4} \\
t_{1} & t_{0} & t_{1} & t_{2} & t_{3} \\
t_{2} & t_{1} & t_{0} & t_{1} & t_{2} \\
t_{3} & t_{2} & t_{1} & t_{0} & t_{1} \\
t_{4} & t_{3} & t_{2} & t_{1} & t_{0}
\end{array}\right] \\
& =\left[\begin{array}{lllll}
t_{4} & t_{3} & t_{2} & t_{1} & t_{0} \\
t_{3} & t_{2} & t_{1} & t_{0} & t_{1} \\
t_{2} & t_{1} & t_{0} & t_{1} & t_{2} \\
t_{1} & t_{0} & t_{1} & t_{2} & t_{3} \\
t_{0} & t_{1} & t_{2} & t_{3} & t_{4}
\end{array}\right]
\end{aligned}
$$

Multiplying a matrix on the right by $J_{n}$ reverses the columns of the matrix, so

$$
\begin{aligned}
J_{5} T_{5} J_{5} & =\left(J_{5} T_{5}\right) J_{5} \\
& =\left[\begin{array}{lllll}
t_{4} & t_{3} & t_{2} & t_{1} & t_{0} \\
t_{3} & t_{2} & t_{1} & t_{0} & t_{1} \\
t_{2} & t_{1} & t_{0} & t_{1} & t_{2} \\
t_{1} & t_{0} & t_{1} & t_{2} & t_{3} \\
t_{0} & t_{1} & t_{2} & t_{3} & t_{4}
\end{array}\right]\left[\begin{array}{lllll}
0 & 0 & 0 & 0 & 1 \\
0 & 0 & 0 & 1 & 0 \\
0 & 0 & 1 & 0 & 0 \\
0 & 1 & 0 & 0 & 0 \\
1 & 0 & 0 & 0 & 0
\end{array}\right] \\
& =\left[\begin{array}{lllll}
t_{0} & t_{1} & t_{2} & t_{3} & t_{4} \\
t_{1} & t_{0} & t_{1} & t_{2} & t_{3} \\
t_{2} & t_{1} & t_{0} & t_{1} & t_{2} \\
t_{3} & t_{2} & t_{1} & t_{0} & t_{1} \\
t_{4} & t_{3} & t_{2} & t_{1} & t_{0}
\end{array}\right]=T_{5}
\end{aligned}
$$

In general, $J_{n} T_{n} J_{n}=T_{n}$. What does this tell us about the eigenvectors of $T_{n}$ ?

First, suppose $\lambda$ is an eigenvalue of $T_{n}$ with multiplicity one, so the associated eigenspace is one-dimensional, and suppose that $T_{n} x=\lambda x$ with $x \neq 0$. Then $J_{n} T_{n} x=$ $\lambda J_{n} x$. Since $J_{n} J_{n}=I_{n}$, it follows that $\left(J_{n} T_{n} J_{n}\right)\left(J_{n} x\right)=$ $\lambda J_{n} x$. Therefore $T_{n}\left(J_{n} x\right)=\lambda\left(J_{n} x\right)$, because $J_{n} T_{n} J_{n}=$ $T_{n}$. Since the $\lambda$-eigenspace of $T_{n}$ is one-dimensonal, it follows that $J_{n} x=c x$ for some constant $c$. Therefore, since $\left\|J_{n} x\right\|=\|x\|$ (that is, $x$ and $J_{n} x$ have the same length), it follows that $c= \pm 1$; that is, $x$ is either symmetric or skew-symmetric. The situation is more complicated if $\lambda$ is a repeated eigenvalue of $T_{n}$ with multiplicity $k$. However, it can be shown that if $k=2 \ell$ then the $\lambda$-eigenspace of $T_{n}$ has a basis consisting of $\ell$ symmetric and $\ell$ skew-symmetric vectors, while if $k=2 \ell+1$ then the $\lambda$-eigenspace has a basis consistng of either $\ell$ symmetric and $\ell+1$ skewsymmetric eigenvectors or $\ell+1$ skew-symmetric and $\ell$ symmetric eigenvectors. In any case, if $n=2 k$ then $T_{n}$ has $k$ symmetric linearly independent eigenvectors and $k$ linearly independent skew-symmetric eigenvectors, while if $n=2 k+1$ then $T_{n}$ has $k+1$ linearly independent symmetric eigenvectors and $k$ linearly independent skewsymmetric vectors.

Recall that the components of the eigenvectors of $T_{n}$ satisfy

$$
\sum_{\ell=-d}^{d} t_{|\ell|} x_{\ell+r}=\lambda x_{r}, \quad 0 \leq r \leq n-1
$$

subject to

$$
\begin{equation*}
x_{r}=0, \quad-d \leq r \leq-1, \quad n \leq r \leq n+d-1 \tag{BC}
\end{equation*}
$$

and are therefore of the form
$x_{r}=\sum_{s=1}^{d}\left(a_{s} z_{s}^{r}(\lambda)+b_{s} z_{s}^{-r}(\lambda)\right), \quad-d \leq r \leq n+d-1$.

However, we now know that we can assume at the outset that the eigenvectors of $T_{n}$ are either symmetric, which means that $x_{n-r+1}=x_{r}$, or skew-symmetric, which means that $x_{n-r+1}=-x_{r}$. So let's build this into (A) at the start!

For clarity
$x_{r}=\sum_{s=1}^{d}\left(a_{s} z_{s}^{r}(\lambda)+b_{s} z_{s}^{-r}(\lambda)\right), \quad-d \leq r \leq n+d-1$.
To get a symmetric eigenvector, let
$b_{s}=a_{s} z_{s}^{n+1}(\lambda) \quad$ so $\quad x_{r}=\sum_{s=1}^{d} a_{s}\left(z_{s}^{r}(\lambda)+z_{s}^{n-r+1}(\lambda)\right)$.
Since $n-(n-r+1)-1=r, x_{n-r+1}=x_{r}$. As for the boundary conditions

$$
\begin{equation*}
x_{r}=0, \quad-d \leq r \leq-1, \quad n \leq r \leq n+d-1 \tag{BC}
\end{equation*}
$$

it is enough to require that $x_{r}=0$ for $-d \leq r \leq-1$, since this and the equality $x_{n-r+1}=x_{r}$ implies that $x_{r}=0$ for $n \leq r \leq n+d-1$.

Therefore, $\lambda$ is an eigenvalue with an associated symmetric eigenvector if and only if

$$
\operatorname{det}\left(\left[z_{s}^{-r}(\lambda)+z_{s}^{n+r+1}(\lambda)\right]_{r, s=1}^{d}\right)=0
$$

For clarity
$x_{r}=\sum_{s=1}^{d}\left(a_{s} z_{s}^{r}(\lambda)+b_{s} z_{s}^{-r}(\lambda)\right), \quad-d \leq r \leq n+d-1$.
To get a skew-symmetric eigenvector, let
$b_{s}=-a_{s} z_{s}^{n+1}(\lambda) \quad$ so $\quad x_{r}=\sum_{s=1}^{d} a_{s}\left(z_{s}^{r}(\lambda)-z_{s}^{n-r+1}(\lambda)\right)$.

Therefore, $\lambda$ is an eigenvalue with an associated skew-symmetric eigenvector if and only if

$$
\operatorname{det}\left(\left[z_{s}^{-r}(\lambda)-z_{s}^{n+r+1}(\lambda)\right]_{r, s=1}^{d}\right)=0
$$

